

# Higher – order Rodrigues dual vectors. Kinematic equations and tangent operator

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The special Euclidean displacement Lie group  $SE_3$  is isomorphic with the Lie group of orthogonal dual tensors  $S\mathbb{O}_3$  [4], [5]. Forwards, we will denote with  $\mathbb{R}$  the set of dual numbers, with  $\underline{V}_3$  the set of dual vectors, with  $\underline{L}(\underline{V}_3, \underline{V}_3)$  the set dual tensors. We will write a generic element from  $S\mathbb{O}_3 = \{ \underline{R} \in \underline{L}(\underline{V}_3, \underline{V}_3) | \underline{R}\underline{R}^T = \underline{I}, \det \underline{R} = 1 \}$  as:

$$\underline{R} = (\underline{I} + \varepsilon \tilde{\rho}) \underline{R}, \varepsilon^2 = 0, \varepsilon \neq 0 \tag{1}$$

where  $\rho$  is the translation vector of the rigid displacement ( $\tilde{\rho}$  denotes the skew-symmetric dual tensor associated to the dual vector  $\rho$ ) and  $\underline{R} \in S\mathbb{O}_3$  the rotation tensor. Also we will write:

$$\underline{R}(\underline{\alpha}, \underline{u}) = \underline{I} + \sin \underline{\alpha} \tilde{u} + (1 - \cos \underline{\alpha}) \tilde{u}^2 = \exp(\underline{\alpha} \tilde{u}) \tag{2}$$

where the unit dual vector  $\underline{u}$  designate the screw-axis of the rigid displacement and  $\underline{\alpha} = \alpha + \varepsilon d$  is the corresponding dual angle [1], [3], [6]. As it was proved in the previous works of the author [5], using the n-th order Cayley transform:

$$cay_n(\cdot) : \underline{V}_3 \rightarrow S\mathbb{O}_3, cay_n(\underline{v}) = (\underline{I} + \tilde{v})^n (\underline{I} - \tilde{v})^{-n}, \forall n \in \mathbb{N}^* \tag{3}$$

a minimal vectorial parameterization for  $SE_3$  can be given. The inverse transformation of  $cay_n$  is a multifunction with n branches given by the below equation [5]:

$$cay_n^{-1}(\cdot) : S\mathbb{O}_3 \rightarrow \underline{V}_3, cay_n^{-1}[\underline{R}(\underline{\alpha}, \underline{u})] = \left( \tan \frac{\underline{\alpha} + 2k\pi}{2n} \right) \underline{u}, k = \overline{0, n-1} \tag{4}$$

The vectors denoted with  $\underline{v}_k = \left( \tan \frac{\underline{\alpha} + 2k\pi}{2n} \right) \underline{u}$ , with  $k = \overline{0, n-1}$  represents a minimal parameterization of the rigid displacement and will be named **higher-order Rodrigues vectors**. If  $\underline{v}$  is one of the Rodrigues dual vectors, then the dual orthogonal tensors that corresponds to the rigid displacement can be computed using the below equation [5].

$$\underline{R} = \underline{I} + \frac{2p_n(|\underline{v}|)q_n(\underline{v})}{(1+|\underline{v}|^2)^n} \tilde{v} + \frac{2q_n^2(|\underline{v}|)}{(1+|\underline{v}|^2)^n} \tilde{v}^2 \tag{5}$$

where  $p_n$  and  $q_n$  are the polynomials given by the following equations:

$$p_n(X) = \sum_{k=0}^{\lfloor n/2 \rfloor} (-1)^k \binom{2k}{n} X^{2k} \tag{6}$$

$$q_n(X) = \sum_{k=0}^{\lfloor (n-1)/2 \rfloor} (-1)^k \binom{2k+1}{n} X^{2k} \tag{7}$$

In the case of a rigid body motion given by a curve  $\underline{R} = \underline{R}(t) \in S\mathbb{O}_3, \forall t \in \mathbb{R}$ , which poses the problem of Rodrigues vector recovery, knowing the dual angular velocity vector (dual twist) in space  $\underline{\omega}$  or body frame  $\underline{\omega}^B$ , we have the following result:

**Theorem:** There are  $\underline{\mathbf{S}}$  and  $\underline{\mathbf{T}}$  two dual non-singular tensors so as:

$$\begin{cases} \underline{\dot{\mathbf{v}}} = \underline{\mathbf{S}}\underline{\boldsymbol{\omega}} \\ \underline{\dot{\mathbf{v}}} = \underline{\mathbf{S}}^T \underline{\boldsymbol{\omega}}^B \end{cases} \quad (8)$$

$$\begin{cases} \underline{\boldsymbol{\omega}} = \underline{\mathbf{T}}\underline{\dot{\mathbf{v}}} \\ \underline{\boldsymbol{\omega}}^B = \underline{\mathbf{T}}^T \underline{\dot{\mathbf{v}}} \end{cases} \quad (9)$$

The following equations will give the tensors  $\underline{\mathbf{S}}$  and  $\underline{\mathbf{T}}$ :

$$\begin{aligned} \underline{\mathbf{S}} &= \frac{p_n(|\underline{\mathbf{v}}|)}{2q_n(\underline{\mathbf{v}})} \underline{\mathbf{I}} - \frac{1}{2} \underline{\tilde{\mathbf{v}}} + \frac{(1+|\underline{\mathbf{v}}|^2)q_n(|\underline{\mathbf{v}}|) - n p_n(|\underline{\mathbf{v}}|)}{2n|\underline{\mathbf{v}}|^2 q_n(\underline{\mathbf{v}})} \underline{\mathbf{v}} \otimes \underline{\mathbf{v}} \\ \underline{\mathbf{T}} &= \frac{2q_n(|\underline{\mathbf{v}}|)p_n(|\underline{\mathbf{v}}|)}{(1+|\underline{\mathbf{v}}|^2)^n} \underline{\mathbf{I}} + \frac{2q_n^2(|\underline{\mathbf{v}}|)}{(1+|\underline{\mathbf{v}}|^2)^n} \underline{\tilde{\mathbf{v}}} + 2 \frac{n(1+|\underline{\mathbf{v}}|^2)^{n-1} - q_n(|\underline{\mathbf{v}}|)p_n(|\underline{\mathbf{v}}|)}{|\underline{\mathbf{v}}|^2(1+|\underline{\mathbf{v}}|^2)^n} \underline{\mathbf{v}} \otimes \underline{\mathbf{v}} \end{aligned} \quad (10)$$

In Eq. (10), we've noted the dual tensor product of two vectors with  $\underline{\mathbf{v}} \otimes \underline{\mathbf{v}}$ .

From the **Rodrigues dual vector parametrization**  $\underline{\mathbf{v}} = \text{tg} \frac{\alpha}{2} \underline{\mathbf{u}}$ ,  $n = 1$ , we will have

$$\begin{aligned} \underline{\mathbf{S}} &= \frac{1}{2} \underline{\mathbf{I}} - \frac{1}{2} \underline{\tilde{\mathbf{v}}} + \frac{1}{2} \underline{\mathbf{v}} \otimes \underline{\mathbf{v}} \\ \underline{\mathbf{T}} &= \frac{2}{1+|\underline{\mathbf{v}}|^2} [\underline{\mathbf{I}} + \underline{\tilde{\mathbf{v}}}] \end{aligned} \quad (11)$$

From the **Wiener-Milenkovic dual vector parametrization**  $\underline{\mathbf{v}} = \text{tg} \frac{\alpha}{4} \underline{\mathbf{u}}$  ( $\underline{\mathbf{v}}_S = -\text{ctg} \frac{\alpha}{4} \underline{\mathbf{u}}$ ),  $n = 2$ , we will obtain:

$$\begin{aligned} \underline{\mathbf{S}} &= \frac{1-|\underline{\mathbf{v}}|^2}{2} \underline{\mathbf{I}} - \frac{1}{2} \underline{\tilde{\mathbf{v}}} + \frac{1}{2} \underline{\mathbf{v}} \otimes \underline{\mathbf{v}} \\ \underline{\mathbf{T}} &= \frac{4}{(1+|\underline{\mathbf{v}}|^2)^2} \left[ (1 - |\underline{\mathbf{v}}|^2) \underline{\mathbf{I}} + 2\underline{\tilde{\mathbf{v}}} + 2\underline{\mathbf{v}} \otimes \underline{\mathbf{v}} \right] \end{aligned} \quad (12)$$

## References

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